Abstract
This project used Python to simulate the effects of properties of Schwarzschild and Kerr black holes, on the paths taken by null geodesics. like Euler's method of integration to Python's scipy library and solve-ivp module, we started with a general simulation of a Schwarzschild black hole. Using the Runge-Kutta 4(5) method. we implemented code that solved four ordinary differential equations, describing the coordinate stein's field equations of relativity [2]. In order
 to validate the underlying physics of our simulation, we tested the code on the precission of Mercury's perihelion, wherein the orbit shifted due to effects captured by general relativity.

## Background

Understanding general relativity and the motion of photons orbiting the accretion disk of black Understanding general relativity and the motion of photons orbiting the accretion disk of black
holes was key for this proiect. The first to do this was Viergutz, who simulated "the shape of accretion disks around Kerr black holes" in 1993 [3]. For a more general expression for black holes in space-time, Kerr metrics have been used for spinning black holes and the same metrics have been constrained for Schwarzschild black holes [4].
$d s^{2}=-\left(1-\frac{r_{s} r}{\sum}\right) c^{2} d t^{2}+\frac{\sum}{\Delta} d r^{2}+\sum d \theta^{2}+\left(r^{2}+a^{2}+\frac{r_{s} r a^{2}}{\sum} \sin ^{2} \theta\right) \sin ^{2} \theta d \phi^{2}-\frac{2 r_{s} r a \sin ^{2} \theta}{\sum} c d t d \phi$ (1)
The differential equations below are for rate-of-change quantities in the radial, theta, and phi directions (in our simulation for Schwarzschild black holes, the affine parameter q was set equal
to one such that $d^{2} t / d q^{2}$ remained constant and equal to one) [4]. For brevity, $w=1-r_{s} r$. where $r_{s}$ is the Schwarzschild radius; and $v=1 / w$. The equations for Kerr black holes include an additional parameter a, which encodes information about the spin ( $a=J / m$, or the total angular momentum per unit mass).

$$
\dot{\varphi}=\frac{2 a r k \sin ^{2} \theta+\left(r^{2}+a^{2} \cos ^{2} \theta-2 r\right) h}{\left(r^{2}+a^{2}\right)\left(r^{2}+a^{2} \cos ^{2} \theta-2 r\right) \sin ^{2} \theta+2 a^{2} r \sin ^{4} \theta}
$$

$\dot{\theta}=\frac{Q+(k a \cos \theta-h \cot \theta)(k a \cos \theta+h \cot \theta)}{\rho^{4}}$
$\dot{r}=\frac{\Delta}{\rho^{2}}\left[k \dot{t}-h \dot{\varphi}-\rho^{2} \dot{\theta}^{2}\right]$

Runge-Kutta 4 Method


The Runge-Kutta 4(5) Method is widely regarded as the most accurate numerical method of integration. It is an algorithm that expands on the classic fourth-order method, and implements an error estimator of order five. For a first-order ODE, each subsequent value for the underlying
function is determined by the present value in addition to a term characterized by: the product of a) the step size, and b) a slope estimated by a function on the right-hand side of the differential equation in question [5]. The higher-order nature of the family of Runge-Kutta method compared to Euler's method to yield less error makes the former an ideal choice.

Schwarzschild and Kerr Simulations
The figure below showcases a stationary black hole with Schwarzschild geodesics programmed into Python. We set six different variables as our state with respect to time rr, theta, phi, and their respective derivatives), and then simulated our values using three geodesics equations with
respect to the derivatives of $r$, theta, and phi. As we initially used spherical coordinates in the respect to the derivatives of $r$, theta, and phi. As we initially used spherical coordinates in the
geodesic equations, we converted them to Cartesian coordinates. Additionally, the initial states were scaled via dividing by an un-normalized speed given by

$$
\sqrt{\dot{r}^{2}+\dot{\theta}^{2}+\dot{\phi}^{2} \sin ^{2} \theta}
$$

We then graphed the equations with sample values for each state using the built-in Solve-IVP function.


For spinning, Kerr black holes (below Figure), we used the same method, but using $r, \theta$ and $\phi$ as our states. In addition to those states, we also had to define a multitude of other constants such as Carter's constant, as well as constants of motion including energy and angular momentum, which resulted from our choice in quantities for the initial state of the particle. It was also crucial to use a different set of 3 main geodesics equations to map our particle. After implementing the required constants and different geodesics equations, we graphed the path of a particle using Solve-IVP as well.



## Mercury Perihelion

To validate the intuition behind the simulation, the program on Mercury's perihelion effect predicted by general relativity has been tested. Due to gravitational forces from other planets, the major axis of Mercury's orbit rotates about the Sun, causing a shift in the line that connects the
Sun to the perihelion of the orbit. Einstein, through his theory of general relativity, called the Sun to the perinelion of the orbit. Einstein, through his theory of general relativity, called the precise prediction of perihelion shift the most critical test of his theory.


Future Work


Figure 5. Reissner-Nordstrom Black Hole [6]
E Event Horizon Telescope on the M8 by the Event Horiz
further validation.

In the future, this project could be extended to Reissner-Nordstrom black holes and KerrNewman black holes, two types of charged Newman black holes, two types of charged
black holes without and with spin, respectively. Thus, the simulations would include all of the types of black holes that were found by solving Einstein's equations of general relativity. Encoding models for the Kerr-Newman black hole would result in a generalization of all possible black hole types. With specific regards to the
simulation, further actions could be undertaken with regards to parallelizing the code to make it more efficient among multiple CPU processors. Rendering our simulations and comparing sors. Rendering odurt to images made possible

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